465.

NOTE ON THE LUNAR THEORY.

[From the Monthly Notices of the Royal Astronomical Society, vol. xxv. (1864—1865), pp. 182—189.]

I ATTEND, in the expressions for the lunar coordinates, only to the coefficients independent of m. Plana's values, taken to the fourth order only, are as follows; for greater simplicity I write a=1; and, instead of nt+ constant, cnt+ constant, gnt+ constant, I write l, c, g respectively; viz., l is the mean longitude, c the mean anomaly, g the mean distance from node: this being so, then r, v, y, denoting the radius vector longitude and latitude respectively, we have

$$\frac{1}{r}(\text{Plana}) = \\ & e - \frac{1}{8} e^3 - \frac{1}{4} \gamma^2 e \quad \cos \quad c \\ & + e^2 - \frac{1}{3} e^4 - \frac{1}{2} \gamma^2 e^2 \quad , \quad 2c \\ & + \frac{9}{8} e^3 \quad , \quad 3c \\ & + \frac{4}{3} e^4 \quad , \quad 4c \\ & - \frac{5}{4} \gamma^2 e^2 \quad , \quad 2g \\ & - \frac{5}{8} \gamma^2 e \quad , \quad c - 2g \\ \\ (\text{but I omit Plana's term} \quad + \frac{1}{8} \gamma^2 e^2 \quad \cos \quad 2c + 2g \text{ which should be} = 0). \\ & v \text{ (Plana)} = l + \\ & + 2 e - \frac{1}{4} e^3 - \frac{1}{2} \gamma^2 e \quad \sin \quad c \\ & + \frac{5}{4} e^2 - \frac{11}{24} e^4 - \frac{15}{16} \gamma^2 e^2 \quad , \quad 2c \\ & + \frac{13}{12} e^3 \quad , \quad 3c \\ & + \frac{103}{96} e^4 \quad , \quad 4c \\ & - \frac{1}{4} \gamma^2 - \frac{9}{16} \gamma^2 e^2 + \frac{1}{8} \gamma^4 \quad , \quad 2g \\ \end{cases}$$

46

49.

To compare these with the elliptic values, it is necessary to write $e(1+\frac{1}{4}\gamma^2)$ in place of e. Making this change, or say reducing Plana's (e, γ) to the elliptic (e, γ) , I write down in a first column the transformed coefficients, and in a second column the elliptic coefficients, as follows:

Plana, with Elliptic e, \gamma

 $+\frac{1}{32} \gamma^4$

Elliptic

where, for greater clearness, I remark that the values called "elliptic" of e, γ , c, g, refer to an ellipse, such that the longitude of the node, and the longitude (in orbit) of the pericentre, vary uniformly with the time,—viz., we have mean distance = 1, excentricity = e, tangent of inclination = γ , mean longitude = l, mean anomaly = c, distance from node = g.

We have therefore

$$\delta \frac{1}{r} = -\frac{5}{4} \gamma^{3} e^{2} \qquad \cos \qquad 2g$$

$$-\frac{5}{8} \gamma^{2} e \qquad , \qquad c - 2g$$

$$\delta v = -\frac{5}{16} \gamma^{2} e^{2} \qquad \sin \qquad 2c$$

$$-\frac{25}{16} \gamma^{2} e^{2} \qquad , \qquad c - 2g$$

$$+\frac{5}{4} \gamma^{2} e \qquad , \qquad c - 2g$$

$$-\frac{5}{16} \gamma^{2} e^{2} \qquad , \qquad c - 2g$$

$$-\frac{5}{16} \gamma^{2} e^{2} \qquad , \qquad 2c - 2g$$

$$\delta y = -\frac{5}{8} \gamma e^{3} + \frac{5}{8} \gamma^{3} e \qquad , \qquad c - g$$

$$+\frac{5}{8} \gamma e^{3} \qquad , \qquad 3c - g$$

$$+\frac{5}{8} \gamma^{3} e \qquad , \qquad c - 3g,$$

viz., these are the increments to be added to the elliptic values of $\frac{1}{r}$, v, y, respectively, in order to obtain the disturbed values of $\frac{1}{r}$, v, y, attending only to the coefficients independent of m; they represent, in fact, the lunar inequalities which rise two orders by integration.

The elliptic values of $\frac{1}{r}$ and y are functions, and that of v, is equal l+, a function, of e, γ , c, g, and the foregoing disturbed values may be obtained by affecting each of 46-2

the quantities e, γ , c, g, and l, with an inequality depending on the argument 2c-2g, viz., these inequalities are

$$\delta e = -\frac{5}{8} \gamma^2 e \cos 2c - 2g$$

$$\delta c = \frac{5}{8} \gamma^2 \sin 2c - 2g$$

$$\delta \gamma = \frac{5}{8} \gamma e^2 \cos 2c - 2g$$

$$\delta g = \frac{5}{8} e^2 \sin 2c - 2g$$

$$\delta l = -\frac{5}{16} \gamma^2 e^2 \sin 2c - 2g.$$

The verification may be effected without difficulty; thus, for instance, starting from the elliptic value of $\frac{1}{n}$, we have to the fourth order

$$\begin{split} \delta \frac{1}{r} &= \delta \begin{pmatrix} e^2 \cos c \\ + e \cos 2c \end{pmatrix} = \begin{pmatrix} -e \sin c \\ -2e^2 \sin 2c \end{pmatrix} \delta c + \begin{pmatrix} \cos c \\ +2e \cos 2c \end{pmatrix} \delta e \\ &= \frac{5}{8} \gamma^2 e \ (-\sin c \sin 2c - 2g - \cos c \cos 2c - 2g) \\ &+ \frac{5}{4} \gamma^2 e^2 (-\sin 2c \sin 2c - 2g - \cos 2c \cos 2c - 2g) \\ &= -\frac{5}{8} \gamma^2 e \cos c - g \\ &- \frac{5}{4} \gamma^2 e^2 \cos 2g, \end{split}$$

which is right; and the verification of the values of δv , δy , may be effected in a similar manner.

I have, in order to fix the ideas, preferred to give in the first instance the foregoing à posteriori proof; but I now inquire generally as to the form of the values of $\frac{1}{r}$, v, y, or say of r, v, y, taking account only of coefficients independent of m; and I proceed to show that these may be obtained from the elliptic values expressed as above in terms of l, e, γ , c, g, by affecting l, e, γ , c, g, each with an inequality depending on the multiple sines or cosines of c-g.

Writing for greater simplicity n=1, we have l=t+L, c=ct+C, g=gt+G, where $c=1-\frac{3}{4}m^2+\&c$, $g=1+\frac{3}{4}m^2+\&c$; viz., c, g, are constants which differ from unity by terms involving m^2 .

The required values of r, v, y, satisfy the *undisturbed* equations of motion, if after the differentiations we write in the coefficients (which coefficients are functions of m through c, g) m=0; that is, if we write in the coefficients c=1, g=1. In fact, the required values of r, v, y, are what the complete values become, upon writing in the coefficients of the complete values m=0; that is, the required values of r, v, y, differ from the complete values by terms the coefficients whereof contain m as a factor; and the disturbed equations differ from the undisturbed equations in that they contain the differential coefficients of the disturbing function; that is, terms the coefficients whereof have the factor m^2 . Imagine the complete values of r, v, y, substituted in the disturbed equations of motion; the resulting equations are satisfied identically; and, therefore, whatever be the value of m; that is, they are satisfied if in these equations respectively

we write m=0: it requires a little consideration to see that this is so, if in the coefficients only we write m=0; but recollecting that c, g, stand for functions ct+C, gt+G, so that, for example, c-g, =(c-g)t+C-G, upon writing therein m=0, becomes equal, not to zero, but to the constant value C-G, the identity subsists in regard to the coefficient of the sine or cosine of each separate argument $\alpha c+\beta g$, and, consequently, it subsists notwithstanding that in the arguments c and g, instead of being each put g are left indeterminate. And granting this (viz. that the equations are satisfied if in the coefficients only we write g be undisturbed equations of motion, if after the differentiations we write in the coefficients g and g are g and g are g being each put g and g are satisfied in the coefficients only we write g be undisturbed equations of motion, if after the differentiations we write in the coefficients g and g are g and g are g are g and g are g and g are g are g and g are g and g are g are g and g are g and g are g are g and g are g are g and g are g are g are g and g are g are g and g are g and g are g are g are g and g are g and g are g and g are g and g are g are g and g are g are g and g are g and g are g and g are g are g are g and g are g are g are g and g are g are g and g are g are g and g are g are g are g and g are g are g are g are g and g are g are g and g are g are g and g are g and g are g are g and g are g are g are g ar

The required values of r, v, y, are of the form $r = \phi(c, g)$, $y = \psi(c, g)$, $v = l + \chi(c, g)$, but writing w = v + c - l, $= c + \chi(c, g)$, the last mentioned property will equally subsist in regard to the functions r, w, y: in fact, v enters into the differential equations only through its differential coefficient $\frac{dv}{dt}$, and the differential coefficients of v and w, that is, of $l + \chi(c, g)$ and $c + \chi(c, g)$, differ only by the quantity c - 1, which becomes c = 0, in virtue of the assumed relations c = 1, c = 1.

Hence the undisturbed equations are satisfied by the values $r = \phi(c, g)$, $y = \psi(c, g)$, $w = c + \chi(c, g)$, when after the differentiations we write in the coefficients c = 1, g = 1; the foregoing values contain t through the quantities c, g, only; and we have, therefore, $\frac{d}{dt} = c \frac{d}{dc} + g \frac{d}{dg}.$

Hence, writing in the coefficients c=1, g=1, we have $\frac{d}{dt}=\frac{d}{dc}+\frac{d}{dg}$; that is, the values $r=\phi$ (c, g), $y=\psi$ (c, g), $w=\chi(c, g)$, regarding r, v, y, as functions of c, g, satisfy the partial differential equations obtained from the undisturbed equations of motion by writing therein $\frac{d}{dc}+\frac{d}{dg}$ in place of $\frac{d}{dt}$. Hence also, considering r, w, y, as functions of c and c-g, then observing that $\left(\frac{d}{dc}+\frac{d}{dg}\right)(c-g)$ is =0, the values of r, v, y, satisfy the partial differential equations obtained by writing $\frac{d}{dc}$ in place of $\frac{d}{dt}$; and inasmuch as these partial differential equations do not contain $\frac{d}{dg}$, they are to be integrated as ordinary differential equations in regard to c as the independent variable, the constants of integration being replaced by arbitrary functions of c-g.

Consider the pure elliptic values of r, v, y, in an elliptic orbit with the following elements: A, the mean distance; N, the mean motion $(N^2A^3=1)$ and therefore $A=N^{-\frac{2}{3}}$; E, the excentricity; Nt+D, the mean anomaly; Nt+H, the mean distance from node; Nt+K, the mean longitude; then writing c in place of t, we have

$$r = N^{-\frac{2}{3}} \operatorname{elqr}(E, Nc + D),$$

$$v(=l-c+w) = l-c+Nc+K+P(E, \Gamma, Nc+D, Nc+H),$$

$$y = Q(E, \Gamma, Nc+D, Nc+H),$$

where N, E, Γ , D, H, K, are arbitrary functions of c-g: P and Q denote given functional expressions. But, in order that r, v, y, considered as functions of c and g may be of the proper form, it is necessary as regards N to write simply N=1; we have then

$$r = \text{elqr}(E, c + D),$$

 $v = l + K + P(E, \Gamma, c + D, c + H),$
 $y = Q(E, \Gamma, c + D, c + H),$

where E, Γ , D, H, K, are arbitrary functions of c-g; or, what is the same thing, writing for these quantities respectively $e + \delta e$, $\gamma + \delta \gamma$, δc , $g - c + \delta g$, δl , where δe , $\delta \gamma$, δc , δg , δl are arbitrary functions of c-g, we have

$$\begin{split} r &= \operatorname{elqr}\left(e + \delta e, \ c + \delta c\right), \\ v &= l + \delta l + P\left(e + \delta e, \ \gamma + \delta \gamma, \ c + \delta c, \ g + \delta g\right), \\ y &= Q\left(e + \delta e, \ \gamma + \delta \gamma, \ c + \delta c, \ g + \delta g\right), \end{split}$$

that is, the values of r, v, y, are obtained from the elliptic values

$$r = \operatorname{elqr}(e, c),$$

 $v = l + P(e, \gamma, c, g),$
 $y = Q(e, \gamma, c, g),$

by affecting each of the quantities e, γ , c, g, l, with an inequality which is a function of c-g.