

XVIII.

THE AUXILIARY FUNCTION T FOR TWO THIN LENSES
CLOSE TOGETHER IN VACUO, AND FOR A SINGLE THIN
LENS IN VACUO

July, 1833.

[Note Book 28, pp. 29-39 (back).]

For a small spheric cap at origin, with curvature = r_i ,

$$z_i = \frac{1}{2} r_i (x_i^2 + y_i^2) + (\frac{1}{2} r_i)^3 (x_i^2 + y_i^2)^2;$$

and

$$\begin{aligned} z_i - p_i x_i - q_i y_i &= r_i^{-1} \{1 - \sqrt{1 + p_i^2 + q_i^2}\} \\ &= -\frac{p_i^2 + q_i^2}{2r_i} + \frac{(p_i^2 + q_i^2)^2}{8r_i}. \end{aligned}$$

Therefore for a series of n spheric refracting or reflecting surfaces close together at origin,*

$$\begin{aligned} T &= -\frac{1}{2} \sum_{(i)1}^n \frac{(\mu_i \alpha_i - \mu_{i-1} \alpha_{i-1})^2 + (\mu_i \beta_i - \mu_{i-1} \beta_{i-1})^2}{r_i (\mu_i \gamma_i - \mu_{i-1} \gamma_{i-1})} \\ &\quad + \frac{1}{8} \sum_{(i)1}^n \frac{\{(\mu_i \alpha_i - \mu_{i-1} \alpha_{i-1})^2 + (\mu_i \beta_i - \mu_{i-1} \beta_{i-1})^2\}^2}{r_i (\mu_i \gamma_i - \mu_{i-1} \gamma_{i-1})^3}; \end{aligned}$$

and

$$V = \mu_n (x_n \alpha_n + y_n \beta_n + z_n \gamma_n) - \mu_0 (x_0 \alpha_0 + y_0 \beta_0 + z_0 \gamma_0) - T.$$

For two thin refracting lenses close together *in vacuo*, with four curved surfaces,†

$$\begin{aligned} V &= x_5 \alpha_4 + y_5 \beta_4 + z_5 \gamma_4 - (x_0 \alpha_0 + y_0 \beta_0 + z_0 \gamma_0) \\ &\quad + \frac{(\mu_1 \alpha_1 - \alpha_0)^2 + (\mu_1 \beta_1 - \beta_0)^2}{2r_1 (\mu_1 \gamma_1 - \gamma_0)} + \frac{(\mu_3 \alpha_3 - \alpha_2)^2 + (\mu_3 \beta_3 - \beta_2)^2}{2r_3 (\mu_3 \gamma_3 - \gamma_2)} \\ &\quad + \frac{(\alpha_2 - \mu_1 \alpha_1)^2 + (\beta_2 - \mu_1 \beta_1)^2}{2r_2 (\gamma_2 - \mu_1 \gamma_1)} + \frac{(\alpha_4 - \mu_3 \alpha_3)^2 + (\beta_4 - \mu_3 \beta_3)^2}{2r_4 (\gamma_4 - \mu_3 \gamma_3)} \\ &\quad - \frac{\{(\mu_1 \alpha_1 - \alpha_0)^2 + (\mu_1 \beta_1 - \beta_0)^2\}^2}{8r_1 (\mu_1 \gamma_1 - \gamma_0)^3} - \frac{\{(\mu_3 \alpha_3 - \alpha_2)^2 + (\mu_3 \beta_3 - \beta_2)^2\}^2}{8r_3 (\mu_3 \gamma_3 - \gamma_2)^3} \\ &\quad - \frac{\{(\alpha_2 - \mu_1 \alpha_1)^2 + (\beta_2 - \mu_1 \beta_1)^2\}^2}{8r_2 (\gamma_2 - \mu_1 \gamma_1)^3} - \frac{\{(\alpha_4 - \mu_3 \alpha_3)^2 + (\beta_4 - \mu_3 \beta_3)^2\}^2}{8r_4 (\gamma_4 - \mu_3 \gamma_3)^3}; \\ \gamma_0 &= \sqrt{1 - \alpha_0^2 - \beta_0^2} = 1 - \frac{1}{2} (\alpha_0^2 + \beta_0^2) - \frac{1}{8} (\alpha_0^2 + \beta_0^2)^2; \\ \gamma_1 &= \&c.; \quad \mu_1 \gamma_1 - \gamma_0 = \mu_1 - 1 - \frac{1}{2} \mu_1 (\alpha_1^2 + \beta_1^2) + \frac{1}{2} (\alpha_0^2 + \beta_0^2); \\ (\mu_1 \gamma_1 - \gamma_0)^{-1} &= (\mu_1 - 1)^{-1} + \frac{\mu_1 (\alpha_1^2 + \beta_1^2) - (\alpha_0^2 + \beta_0^2)}{2(\mu_1 - 1)^2}; \end{aligned}$$

* [The thickness of the lenses, although in general of the order of $\alpha_0^2 + \beta_0^2$, is here neglected. The argument used here is that of (K⁷) of the Third Supplement (p. 216).]† [The point x_5, y_5, z_5 , is any point on the final ray.]

(July 22, 1833) the approximate conditions of the stationary value* with respect to $\alpha_1, \alpha_2, \alpha_3$ are

$$\frac{\mu_1 \alpha_1 - \alpha_0}{r_1} + \frac{\alpha_2 - \mu_1 \alpha_1}{r_2} = 0, \quad \frac{\alpha_2 - \mu_1 \alpha_1}{r_2 (\mu_1 - 1)} + \frac{\mu_3 \alpha_3 - \alpha_2}{r_3 (\mu_3 - 1)} = 0, \quad \frac{\mu_3 \alpha_3 - \alpha_2}{r_3} + \frac{\alpha_4 - \mu_3 \alpha_3}{r_4} = 0;$$

that is,

$$\left. \begin{aligned} \mu_1 (r_1 - r_2) \alpha_1 &= r_1 \alpha_2 - r_2 \alpha_0 \\ \mu_3 (r_3 - r_4) \alpha_3 &= r_3 \alpha_4 - r_4 \alpha_2 \end{aligned} \right\}$$

and

$$\alpha_2 \{r_2 (\mu_1 - 1) - r_3 (\mu_3 - 1)\} = \frac{r_2 (\mu_1 - 1) (r_3 \alpha_4 - r_4 \alpha_2)}{r_3 - r_4} - \frac{r_3 (\mu_3 - 1) (r_1 \alpha_2 - r_2 \alpha_0)}{r_1 - r_2}$$

$$\therefore \dagger \alpha_2 \left\{ r_2 (\mu_1 - 1) - r_3 (\mu_3 - 1) + \frac{r_2 r_4 (\mu_1 - 1)}{r_3 - r_4} + \frac{r_1 r_3 (\mu_3 - 1)}{r_1 - r_2} \right\} = r_2 r_3 \left\{ \frac{\alpha_4 (\mu_1 - 1)}{r_3 - r_4} + \frac{\alpha_0 (\mu_3 - 1)}{r_1 - r_2} \right\}$$

$$\therefore \alpha_2 \left\{ \frac{\mu_1 - 1}{r_3 - r_4} + \frac{\mu_3 - 1}{r_1 - r_2} \right\} = \frac{\alpha_4 (\mu_1 - 1)}{r_3 - r_4} + \frac{\alpha_0 (\mu_3 - 1)}{r_1 - r_2}$$

$$\therefore \alpha_2 \{(\mu_1 - 1) (r_1 - r_2) + (\mu_3 - 1) (r_3 - r_4)\} = \alpha_0 (\mu_3 - 1) (r_3 - r_4) + \alpha_4 (\mu_1 - 1) (r_1 - r_2),$$

\(\therefore\) finally ‡

$$\alpha_2 = \frac{\alpha_0 P_1 + \alpha_4 P_2}{P_1 + P_2} = \frac{\alpha' P_1 + \alpha P_2}{P_1 + P_2}; \quad \beta_2 = \frac{\beta P_1 + \beta' P_2}{P_1 + P_2},$$

using P, P_1 , to denote the powers of the 1st and 2nd lens. That is,

$$P (\alpha_4 - \alpha_2) = P_1 (\alpha_2 - \alpha_0):$$

an equation which probably admits of some simple geometrical enunciation.

$$\mu_1 \alpha_1 = \alpha_2 + \frac{r_2 (\alpha_2 - \alpha_0)}{r_1 - r_2} = \alpha_2 + \frac{(\mu_1 - 1) r_2 (\alpha_2 - \alpha_0)}{P} = \alpha_2 + \frac{(\mu_1 - 1) r_2 (\alpha_4 - \alpha_0)}{P + P_1},$$

$$= \alpha_0 + \frac{r_1 (\alpha_2 - \alpha_0)}{r_1 - r_2} = \alpha_0 + \frac{(\mu_1 - 1) r_1 (\alpha_2 - \alpha_0)}{P} = \alpha_0 + \frac{(\mu_1 - 1) r_1 (\alpha_4 - \alpha_0)}{P + P_1};$$

$$\mu_3 \alpha_3 = \alpha_4 + \frac{r_4 (\alpha_4 - \alpha_2)}{r_3 - r_4} = \alpha_4 + \frac{(\mu_3 - 1) r_4 (\alpha_4 - \alpha_2)}{P_1} = \alpha_4 + \frac{(\mu_3 - 1) r_4 (\alpha_4 - \alpha_0)}{P + P_1},$$

$$= \alpha_2 + \frac{r_3 (\alpha_4 - \alpha_2)}{r_3 - r_4} = \alpha_2 + \frac{(\mu_3 - 1) r_3 (\alpha_4 - \alpha_2)}{P_1} = \alpha_2 + \frac{(\mu_3 - 1) r_3 (\alpha_4 - \alpha_0)}{P + P_1},$$

$$= \frac{\alpha_4 P + \alpha_0 P_1 + (\mu_3 - 1) r_3 (\alpha_4 - \alpha_0)}{P + P_1} = \alpha_4 + \frac{(\mu_3 - 1) r_4 (\alpha_4 - \alpha_0)}{P + P_1}, \text{ as before:}$$

and similarly

$$\mu_1 \beta_1 = \beta_0 + \frac{(\mu_1 - 1) r_1 (\beta_4 - \beta_0)}{P + P_1}; \quad \mu_3 \beta_3 = \beta_4 + \frac{(\mu_3 - 1) r_4 (\beta_4 - \beta_0)}{P + P_1}.$$

* [That is, of T ; the function V plays no essential part in this investigation.]

† [In the manuscripts of this period, Hamilton used this symbol \therefore consistently in the sense of "therefore." The symbol was used in this sense in the seventeenth century by W. Oughtred and J. H. Rahn; see F. Cajori, *History of Mathematical Notations*, vol. I (Chicago, 1928), pp. 190, 211.]

‡ [Hamilton here reverts to his customary notation, in which a', β', γ' refer to the incident ray, a, β, γ to the emergent ray.]

These values of $\alpha_1, \alpha_2, \alpha_3, \beta_1, \beta_2, \beta_3$ are to be substituted in V or T . Thus,*

$$\begin{aligned} -T^{(2)} &= \frac{(\mu_1\alpha_1 - \alpha_0)^2}{2r_1(\mu_1 - 1)} - \frac{(\alpha_2 - \mu_1\alpha_1)^2}{2r_2(\mu_1 - 1)} + \frac{(\mu_3\alpha_3 - \alpha_2)^2}{2r_3(\mu_3 - 1)} - \frac{(\alpha_4 - \mu_3\alpha_3)^2}{2r_4(\mu_3 - 1)} \\ &\quad + \text{same function of the } \beta\text{'s} \\ &= \frac{\frac{1}{2}P \{(\alpha_4 - \alpha_0)^2 + (\beta_4 - \beta_0)^2\}}{(P + P_1)^2} + \frac{\frac{1}{2}P_1 \{(\alpha_4 - \alpha_0)^2 + (\beta_4 - \beta_0)^2\}}{(P + P_1)^2} \\ &= \frac{(\alpha - \alpha')^2 + (\beta - \beta')^2}{2(P + P_1)}. \end{aligned}$$

Also

$$\begin{aligned} \Delta\sigma_0 &= \mu_1\alpha_1 - \alpha_0 = \lambda_0(\alpha - \alpha'), & \Delta\tau_0 &= \lambda_0(\beta - \beta'), & \lambda_0 &= \frac{(\mu_1 - 1)r_1}{P + P_1}; \\ \Delta\sigma_1 &= \alpha_2 - \mu_1\alpha_1 = \lambda_1(\alpha - \alpha'), & \Delta\tau_1 &= \lambda_1(\beta - \beta'), & \lambda_1 &= -\frac{(\mu_1 - 1)r_2}{P + P_1}; \\ \Delta\sigma_2 &= \mu_3\alpha_3 - \alpha_2 = \lambda_2(\alpha - \alpha'), & \Delta\tau_2 &= \lambda_2(\beta - \beta'), & \lambda_2 &= \frac{(\mu_3 - 1)r_3}{P + P_1}; \\ \Delta\sigma_3 &= \alpha_4 - \mu_3\alpha_3 = \lambda_3(\alpha - \alpha'), & \Delta\tau_3 &= \lambda_3(\beta - \beta'), & \lambda_3 &= -\frac{(\mu_3 - 1)r_4}{P + P_1}; \end{aligned}$$

and

$$\begin{aligned} \Delta v_0 &= \mu_1\gamma_1 - \gamma_0 = \mu_1 - 1 - \frac{1}{2}\mu_1(\alpha_1^2 + \beta_1^2) + \frac{1}{2}(\alpha_0^2 + \beta_0^2) \\ &= \mu_1 - 1 - \frac{(\alpha_0 + \Delta\sigma_0)^2 + (\beta_0 + \Delta\tau_0)^2}{2\mu_1} + \frac{\alpha_0^2 + \beta_0^2}{2} \\ &= (\mu_1 - 1) \left\{ 1 + \frac{\alpha_0^2 + \beta_0^2}{2\mu_1} - \frac{\alpha_0\Delta\sigma_0 + \beta_0\Delta\tau_0}{\mu_1(\mu_1 - 1)} - \frac{(\Delta\sigma_0)^2 + (\Delta\tau_0)^2}{2\mu_1(\mu_1 - 1)} \right\}; \\ \Delta v_1 &= \gamma_2 - \mu_1\gamma_1 = 1 - \mu_1 - \frac{1}{2}(\alpha_2^2 + \beta_2^2) + \frac{1}{2}\mu_1(\alpha_1^2 + \beta_1^2) \\ &= -(\mu_1 - 1) \left\{ 1 + \frac{\alpha_2^2 + \beta_2^2}{2(\mu_1 - 1)} - \frac{\mu_1(\alpha_1^2 + \beta_1^2)}{2(\mu_1 - 1)} \right\}; \\ \Delta v_2 &= \mu_3\gamma_3 - \gamma_2 = (\mu_3 - 1) \left\{ 1 - \frac{\mu_3(\alpha_3^2 + \beta_3^2)}{2(\mu_3 - 1)} + \frac{\alpha_2^2 + \beta_2^2}{2(\mu_3 - 1)} \right\}; \\ \Delta v_3 &= \gamma_4 - \mu_3\gamma_3 = -(\mu_3 - 1) \left\{ 1 - \frac{\mu_3(\alpha_3^2 + \beta_3^2)}{2(\mu_3 - 1)} + \frac{\alpha_4^2 + \beta_4^2}{2(\mu_3 - 1)} \right\}; \end{aligned}$$

or better thus: change μ_1, μ_3 to μ, μ ; and put

$$\alpha_4^2 + \beta_4^2 = \epsilon, \quad \alpha_0^2 + \beta_0^2 = \epsilon', \quad \alpha_4\alpha_0 + \beta_4\beta_0 = \epsilon;$$

then

$$\begin{aligned} (\alpha_4 - \alpha_0)^2 + (\beta_4 - \beta_0)^2 &= \epsilon - 2\epsilon + \epsilon'; \\ \Delta\sigma_0^2 + \Delta\tau_0^2 &= \lambda_0^2(\epsilon - 2\epsilon + \epsilon'); & \Delta\sigma_2^2 + \Delta\tau_2^2 &= \lambda_2^2(\epsilon - 2\epsilon + \epsilon'); \\ \Delta\sigma_1^2 + \Delta\tau_1^2 &= \lambda_1^2(\quad); & \Delta\sigma_3^2 + \Delta\tau_3^2 &= \lambda_3^2(\quad); \\ \gamma_0 &= 1 - \frac{1}{2}(\alpha_0^2 + \beta_0^2) = 1 - \frac{1}{2}\epsilon'; & \gamma_4 &= 1 - \frac{1}{2}(\alpha_4^2 + \beta_4^2) = 1 - \frac{1}{2}\epsilon; \end{aligned}$$

* [$T^{(2)}$ denotes the part of the expansion of T which is of the second order in $\alpha, \beta, \alpha', \beta'$.]

$$\mu_1 \gamma_1 = \mu - \frac{\mu}{2} (\alpha_1^2 + \beta_1^2) = \mu - \frac{1}{2\mu} \{ \alpha' + \lambda_0 (\alpha - \alpha')^2 + \beta' + \lambda_0 (\beta - \beta')^2 \}$$

$$= \mu - \frac{1}{2\mu} \{ \epsilon' + 2\lambda_0 (\epsilon, - \epsilon') + \lambda_0^2 (\epsilon - 2\epsilon, + \epsilon') \};$$

$$\gamma_2 = 1 - \frac{1}{2} (\alpha_2^2 + \beta_2^2) = 1 - \frac{P^2 \epsilon + 2PP, \epsilon, + P,^2 \epsilon'}{2(P + P,)^2};$$

$$\mu_3 \gamma_3 = \mu, - \frac{\mu,}{2} (\alpha_3^2 + \beta_3^2) = \mu, - \frac{1}{2\mu,} \{ (\alpha - \lambda_3 \alpha - \alpha')^2 + (\beta - \lambda_3 \beta - \beta')^2 \}$$

$$= \mu, - \frac{1}{2\mu,} \{ \epsilon - 2\lambda_3 (\epsilon - \epsilon,) + \lambda_3^2 (\epsilon - 2\epsilon, + \epsilon') \};$$

$$\therefore \left\{ \begin{aligned} \Delta v_0 &= \mu_1 \gamma_1 - \gamma_0 = (\mu - 1) \left\{ 1 + \frac{\epsilon'}{2\mu} - \frac{\lambda_0 (\epsilon, - \epsilon')}{\mu (\mu - 1)} - \frac{\lambda_0^2 (\epsilon - 2\epsilon, + \epsilon')}{2\mu (\mu - 1)} \right\}; \\ \Delta v_1 &= \gamma_2 - \mu_1 \gamma_1 = -(\mu - 1) \left\{ 1 + \frac{P^2 \epsilon + 2PP, \epsilon, + P,^2 \epsilon'}{2(\mu - 1)(P + P,)^2} \right. \\ &\quad \left. - \frac{\epsilon' + 2\lambda_0 (\epsilon, - \epsilon') + \lambda_0^2 (\epsilon - 2\epsilon, + \epsilon')}{2\mu (\mu - 1)} \right\}; \\ \Delta v_2 &= \mu_3 \gamma_3 - \gamma_2 = (\mu, - 1) \left\{ 1 + \frac{P^2 \epsilon + 2PP, \epsilon, + P,^2 \epsilon'}{2(\mu, - 1)(P + P,)^2} \right. \\ &\quad \left. - \frac{\epsilon - 2\lambda_3 (\epsilon - \epsilon,) + \lambda_3^2 (\epsilon - 2\epsilon, + \epsilon')}{2\mu, (\mu, - 1)} \right\}; \\ \Delta v_3 &= \gamma_4 - \mu_3 \gamma_3 = -(\mu, - 1) \left\{ 1 + \frac{\epsilon}{2\mu,} + \frac{\lambda_3 (\epsilon - \epsilon,)}{\mu, (\mu, - 1)} - \frac{\lambda_3^2 (\epsilon - 2\epsilon, + \epsilon')}{2\mu, (\mu, - 1)} \right\}. \end{aligned} \right.$$

$$T = -\frac{1}{2} \sum_{(i)1}^4 \frac{(\Delta \sigma_{i-1})^2 + (\Delta \tau_{i-1})^2}{r_i \Delta v_{i-1}} + \frac{1}{8} \sum_{(i)1}^4 \frac{[(\Delta \sigma_{i-1})^2 + (\Delta \tau_{i-1})^2]^2}{r_i (\Delta v_{i-1})^3};$$

$$r_1 (\mu - 1) = \lambda_0 (P + P,); \quad -r_4 (\mu, - 1) = \lambda_3 (P + P,);$$

$$-r_2 (\mu - 1) = \lambda_1 (P + P,); \quad r_3 (\mu, - 1) = \lambda_2 (P + P,);$$

$$\Delta v_0 = (\mu - 1)(1 - \xi_0); \quad \Delta v_1 = -(\mu - 1)(1 - \xi_1);$$

$$\Delta v_2 = (\mu, - 1)(1 - \xi_2); \quad \Delta v_3 = -(\mu, - 1)(1 - \xi_3);$$

$$r_1 \Delta v_0 = \lambda_0 (P + P,)(1 - \xi_0); \quad r_2 \Delta v_1 = \lambda_1 (P + P,)(1 - \xi_1);$$

$$r_3 \Delta v_2 = \lambda_2 (P + P,)(1 - \xi_2); \quad r_4 \Delta v_3 = \lambda_3 (P + P,)(1 - \xi_3);$$

therefore observing that $\lambda_0 + \lambda_1 + \lambda_2 + \lambda_3 = 1$, we get as before*

$$T^{(2)} = -\frac{(\epsilon - 2\epsilon, + \epsilon')}{2(P + P,)},$$

and

$$T^{(4)} = -\frac{(\lambda_0 \xi_0 + \lambda_1 \xi_1 + \lambda_2 \xi_2 + \lambda_3 \xi_3)(\epsilon - 2\epsilon, + \epsilon')}{2(P + P,)} \\ + \frac{(\epsilon - 2\epsilon, + \epsilon')^2}{8(P + P,)} \left\{ \frac{\lambda_0^3 + \lambda_1^3}{(\mu - 1)^2} + \frac{\lambda_2^3 + \lambda_3^3}{(\mu, - 1)^2} \right\};$$

* [For a justification of the method of approximation employed, see Appendix, Note 24, p. 507.]

in which

$$\begin{aligned} \frac{\lambda_0^3 + \lambda_1^3}{(\mu - 1)^2} &= \frac{(\mu - 1)(r_1^3 - r_2^3)}{(P + P_1)^3} = \frac{P(r_1^2 + r_1 r_2 + r_2^2)}{(P + P_1)^3}, \\ \frac{\lambda_2^3 + \lambda_3^3}{(\mu_1 - 1)^2} &= \frac{(\mu_1 - 1)(r_3^3 - r_4^3)}{(P + P_1)^3} = \frac{P_1(r_3^2 + r_3 r_4 + r_4^2)}{(P + P_1)^3}; \\ -\lambda_0 \xi_0 &= \frac{\lambda_0 \epsilon'}{2\mu} - \frac{\lambda_0^2(\epsilon - \epsilon')}{\mu(\mu - 1)} - \frac{\lambda_0^3(\epsilon - 2\epsilon_1 + \epsilon')}{2\mu(\mu - 1)}, \\ -\lambda_1 \xi_1 &= -\frac{\lambda_1 \epsilon'}{2\mu(\mu - 1)} - \frac{\lambda_1 \lambda_0(\epsilon - \epsilon')}{\mu(\mu - 1)} - \frac{\lambda_1 \lambda_0^2(\epsilon - 2\epsilon_1 + \epsilon')}{2\mu(\mu - 1)} \\ &\quad + \frac{\lambda_1}{2(\mu - 1)} \{(\lambda_1 + \lambda_0)^2 \epsilon + 2(\lambda_1 + \lambda_0)(\lambda_3 + \lambda_2)\epsilon_1 + (\lambda_3 + \lambda_2)^2 \epsilon'\}, \\ -\lambda_2 \xi_2 &= \frac{\lambda_2}{2(\mu_1 - 1)} \{(\lambda_1 + \lambda_0)^2 \epsilon + \dots\} \\ &\quad - \frac{\lambda_2 \epsilon}{2\mu_1(\mu_1 - 1)} + \frac{\lambda_2 \lambda_3(\epsilon - \epsilon_1)}{\mu_1(\mu_1 - 1)} - \frac{\lambda_2 \lambda_3^2(\epsilon - 2\epsilon_1 + \epsilon')}{2\mu_1(\mu_1 - 1)}, \\ -\lambda_3 \xi_3 &= \frac{\lambda_3 \epsilon}{2\mu_1} + \frac{\lambda_3^2(\epsilon - \epsilon_1)}{\mu_1(\mu_1 - 1)} - \frac{\lambda_3^3(\epsilon - 2\epsilon_1 + \epsilon')}{2\mu_1(\mu_1 - 1)}; \end{aligned}$$

therefore since $\lambda_1 = -\lambda_0 + \frac{P}{P + P_1}$, and $\lambda_2 = -\lambda_3 + \frac{P_1}{P + P_1}$, we have

$$\begin{aligned} -2(\lambda_0 \xi_0 + \lambda_1 \xi_1 + \lambda_2 \xi_2 + \lambda_3 \xi_3) &= \frac{\lambda_0 \epsilon'}{\mu - 1} - \frac{2\lambda_0 P(\epsilon - \epsilon')}{\mu(\mu - 1)(P + P_1)} - \frac{\lambda_0^2 P(\epsilon - 2\epsilon_1 + \epsilon')}{\mu(\mu - 1)(P + P_1)} \\ &\quad - \frac{P \epsilon'}{\mu(\mu - 1)(P + P_1)} + \left\{ \frac{1}{\mu - 1} \left(\frac{P}{P + P_1} - \lambda_0 \right) + \frac{1}{\mu_1 - 1} \left(\frac{P_1}{P + P_1} - \lambda_3 \right) \right\} \frac{P^2 \epsilon + 2PP_1 \epsilon_1 + P_1^2 \epsilon'}{(P + P_1)^2} \\ &\quad - \frac{P_1 \epsilon}{\mu_1(\mu_1 - 1)(P + P_1)} + \frac{\lambda_3 \epsilon}{\mu_1 - 1} + \frac{2\lambda_3 P_1(\epsilon - \epsilon_1)}{\mu_1(\mu_1 - 1)(P + P_1)} - \frac{\lambda_3^2 P_1(\epsilon - 2\epsilon_1 + \epsilon')}{\mu_1(\mu_1 - 1)(P + P_1)}. \end{aligned}$$

If we put

$$T^{(4)} = Q\epsilon^2 + Q_1\epsilon\epsilon_1 + Q_2\epsilon\epsilon_1' + Q_3\epsilon_1'^2 + Q_4\epsilon_1\epsilon_1' + Q_5\epsilon_1'^2,$$

then

$$\begin{aligned} Q &= \frac{P(r_1^2 + r_1 r_2 + r_2^2) + P_1(r_3^2 + r_3 r_4 + r_4^2)}{8(P + P_1)^4} \\ &\quad - \frac{1}{2(P + P_1)} \left\{ \lambda_0 \frac{\delta \xi_0}{\delta \epsilon} + \lambda_1 \frac{\delta \xi_1}{\delta \epsilon} + \lambda_2 \frac{\delta \xi_2}{\delta \epsilon} + \lambda_3 \frac{\delta \xi_3}{\delta \epsilon} \right\}; \\ Q_1 &= -\frac{P(r_1^2 + r_1 r_2 + r_2^2) + P_1(r_3^2 + r_3 r_4 + r_4^2)}{2(P + P_1)^4} \\ &\quad + \frac{1}{P + P_1} \left\{ \lambda_0 \frac{\delta \xi_0}{\delta \epsilon} + \dots \right\} \\ &\quad - \frac{1}{2(P + P_1)} \left\{ \lambda_0 \frac{\delta \xi_0}{\delta \epsilon} + \dots \right\}; \text{ \&c.} \end{aligned}$$

July 23d.) Passing to a single lens, $r_3 = 0$, $r_4 = 0$, therefore $\lambda_2 = 0$, $\lambda_3 = 0$, $\lambda_0 + \lambda_1 = 1$, $P = 0$,

$$-2(\lambda_0 \xi_0 + \lambda_1 \xi_1) = \frac{\lambda_0 \epsilon'}{\mu - 1} - \frac{2\lambda_0(\epsilon - \epsilon')}{\mu(\mu - 1)} - \frac{\lambda_0^2(\epsilon - 2\epsilon + \epsilon')}{\mu(\mu - 1)} - \frac{\epsilon'}{\mu(\mu - 1)} + \frac{\epsilon \lambda_1}{\mu - 1}$$

$$= \frac{r_1}{P\mu} \left\{ \mu \epsilon' - 2(\epsilon - \epsilon') - \frac{(\mu - 1)r_1(\epsilon - 2\epsilon + \epsilon')}{P} - \frac{\epsilon' P}{(\mu - 1)r_1} - \frac{\epsilon r_2 \mu}{r_1} \right\};$$

$$Q = \frac{r_1^2 + r_1 r_2 + r_2^2}{8P^3} - \frac{(\mu - 1)r_1^2}{4\mu P^3} - \frac{r_2}{4P^2};$$

$$Q' = -\frac{r_1^2 + r_1 r_2 + r_2^2}{2P^3} + \frac{(\mu - 1)r_1^2}{2\mu P^3} + \frac{r_2}{2P^2} - \frac{r_1}{2\mu P^2} + \frac{(\mu - 1)r_1^2}{2\mu P^3};$$

and at the same time $r_2 = r_1 - \frac{P}{\mu - 1}$, therefore*

$$Q = \frac{r_1^2}{4P^3} \left(\frac{1}{2} + \frac{1}{\mu} \right) - \frac{r_1}{4P^2} \left(1 + \frac{\frac{3}{2}}{\mu - 1} \right) + \frac{1 + 2(\mu - 1)}{8P(\mu - 1)^2}$$

$$= \frac{(\mu + 2)r_1^2}{8\mu P^3} - \frac{r_1(2\mu + 1)}{8P^2(\mu - 1)} + \frac{2\mu - 1}{8P(\mu - 1)^2};$$

$$Q' = \frac{r_1^2}{P^3} \left(-\frac{3}{2} + \frac{\mu - 1}{\mu} \right) + \frac{r_1}{2P^2} \left(\frac{3}{\mu - 1} + 1 - \frac{1}{\mu} \right) - \frac{\mu}{2P(\mu - 1)^2}$$

$$= -\frac{(\mu + 2)r_1^2}{2\mu P^3} + \frac{r_1(\mu^2 + \mu + 1)}{2P^2\mu(\mu - 1)} - \frac{\mu}{2P(\mu - 1)^2}; \text{ \&c.}$$

July 27th.) For this case of a SINGLE THIN LENS μ , IN VACUO, AT ORIGIN,

$$V = x\alpha + y\beta + z\gamma - x'\alpha' - y'\beta' - z'\gamma' - T^{(2)} - T^{(4)},$$

$$T^{(2)} + T^{(4)} = -\frac{(\mu\alpha_1 - \alpha')^2 + (\mu\beta_1 - \beta')^2}{2r_1(\mu\gamma_1 - \gamma')} + \frac{(\mu\alpha_1 - \alpha)^2 + (\mu\beta_1 - \beta)^2}{2r_2(\mu\gamma_1 - \gamma)}$$

$$+ \frac{\{(\mu\alpha_1 - \alpha')^2 + (\mu\beta_1 - \beta')^2\}^2}{8r_1(\mu - 1)^3} - \frac{\{(\mu\alpha_1 - \alpha)^2 + (\mu\beta_1 - \beta)^2\}^2}{8r_2(\mu - 1)^3};$$

$$-\frac{\mu\alpha_1 + \alpha'}{r_1} + \frac{\mu\alpha_1 - \alpha}{r_2} = 0, \quad \alpha_1 = \frac{\frac{\alpha}{r_2} - \frac{\alpha'}{r_1}}{\mu \left(\frac{1}{r_2} - \frac{1}{r_1} \right)} = \frac{\alpha r_1 - \alpha' r_2}{\mu(r_1 - r_2)};$$

$$\mu\alpha_1 - \alpha' = \frac{(\alpha - \alpha')r_1}{r_1 - r_2}, \quad \mu\alpha_1 - \alpha = \frac{(\alpha - \alpha')r_2}{r_1 - r_2};$$

similarly for the β 's; therefore putting as before

$$\epsilon = \alpha^2 + \beta^2, \quad \epsilon = \alpha\alpha' + \beta\beta', \quad \epsilon' = \alpha'^2 + \beta'^2,$$

* [These expressions have been corrected: in the first line, the MS. reads $-\frac{\frac{3}{2}}{\mu - 1}$ instead of $+\frac{\frac{3}{2}}{\mu - 1}$, and, in the second, $(2\mu - 5)$ instead of $(2\mu + 1)$. The compact method which follows is independent of these results.]

we have

$$\left. \begin{aligned} (\mu\alpha_1 - \alpha')^2 + (\mu\beta_1 - \beta')^2 &= \frac{r_1^2(\epsilon - 2\epsilon + \epsilon')}{(r_1 - r_2)^2} \\ (\mu\alpha_1 - \alpha)^2 + (\mu\beta_1 - \beta)^2 &= \frac{r_2^2(\epsilon - 2\epsilon + \epsilon')}{(r_1 - r_2)^2} \end{aligned} \right\} \therefore T^{(2)} = \frac{-(\epsilon - 2\epsilon + \epsilon')}{2(\mu - 1)(r_1 - r_2)};$$

$$\begin{aligned} &\left\{ T^{(4)} - \frac{(r_1^2 + r_1r_2 + r_2^2)(\epsilon - 2\epsilon + \epsilon')^2}{8(\mu - 1)^3(r_1 - r_2)^3} \right\} \div \frac{\epsilon - 2\epsilon + \epsilon'}{4(\mu - 1)^2(r_1 - r_2)^2} \\ &= 2r_1(\overline{\mu\gamma_1 - \gamma' - \mu - 1}) - 2r_2(\overline{\mu\gamma_1 - \gamma - \mu - 1}) \\ &= r_1\{\epsilon' - \mu(\alpha_1^2 + \beta_1^2)\} - r_2\{\epsilon - \mu(\alpha_1^2 + \beta_1^2)\} \\ &= r_1\epsilon' - r_2\epsilon - \mu(r_1 - r_2)(\alpha_1^2 + \beta_1^2) \\ &= r_1\epsilon' - r_2\epsilon - \frac{\epsilon r_1^2 - 2\epsilon r_1 r_2 + \epsilon' r_2^2}{\mu(r_1 - r_2)}; \end{aligned}$$

$$\begin{aligned} \therefore \frac{8(\mu - 1)^3(r_1 - r_2)^3 T^{(4)}}{\epsilon - 2\epsilon + \epsilon'} &= (r_1^2 + r_1r_2 + r_2^2)(\epsilon - 2\epsilon + \epsilon') + 2(\mu - 1)(r_1 - r_2)(r_1\epsilon' - r_2\epsilon) \\ &\quad - 2(1 - \mu^{-1})(\epsilon r_1^2 - 2\epsilon r_1 r_2 + \epsilon' r_2^2) \\ &= \epsilon\{r_1^2(-1 + 2\mu^{-1}) + r_1r_2(-2\mu + 3) + r_2^2(2\mu - 1)\} - 2\epsilon\{r_1^2 + r_1r_2(-1 + 2\mu^{-1}) + r_2^2\} \\ &\quad + \epsilon'\{r_1^2(2\mu - 1) + r_1r_2(-2\mu + 3) + r_2^2(-1 + 2\mu^{-1})\}; \end{aligned}$$

therefore putting

$$T^{(4)} = \epsilon^2 Q + \epsilon\epsilon' Q_1 + \epsilon\epsilon' Q' + \epsilon_1^2 Q_{11} + \epsilon_1\epsilon' Q_{11}' + \epsilon'^2 Q_{11}';$$

we have*

$$\left. \begin{aligned} Q &= \frac{r_1^2(-\mu + 2) + r_1r_2(-2\mu^2 + 3\mu) + r_2^2(2\mu^2 - \mu)}{8\mu(\mu - 1)^3(r_1 - r_2)^3}; \\ Q_1 &= \frac{r_1^2 + r_1r_2(-\mu^2 + \mu + 1) + r_2^2\mu^2}{-2\mu(\mu - 1)^3(r_1 - r_2)^3}; \\ Q' &= \frac{r_1^2(\mu^2 - \mu + 1) + r_1r_2(-2\mu^2 + 3\mu) + r_2^2(\mu^2 - \mu + 1)}{4\mu(\mu - 1)^3(r_1 - r_2)^3}; \\ Q_{11} &= \frac{r_1^2\mu + r_1r_2(-\mu + 2) + r_2^2\mu}{2\mu(\mu - 1)^3(r_1 - r_2)^3}; \\ Q_{11}' &= \frac{r_1^2\mu^2 + r_1r_2(-\mu^2 + \mu + 1) + r_2^2}{-2\mu(\mu - 1)^3(r_1 - r_2)^3}; \\ Q_{11}'' &= \frac{r_1^2(2\mu^2 - \mu) + r_1r_2(-2\mu^2 + 3\mu) + r_2^2(-\mu + 2)}{8\mu(\mu - 1)^3(r_1 - r_2)^3}. \end{aligned} \right\} \begin{array}{l} \text{Single thin lens} \\ \text{in vacuo.} \end{array}$$

Let p be power; $p = (\mu - 1)(r_1 - r_2)$: then

$$T^{(2)} = -\frac{1}{2p}(\epsilon - 2\epsilon + \epsilon');$$

[The manuscript ends at this point.]

* [The expression for Q'' has been corrected; the MS. reads "3" instead of "3μ."]